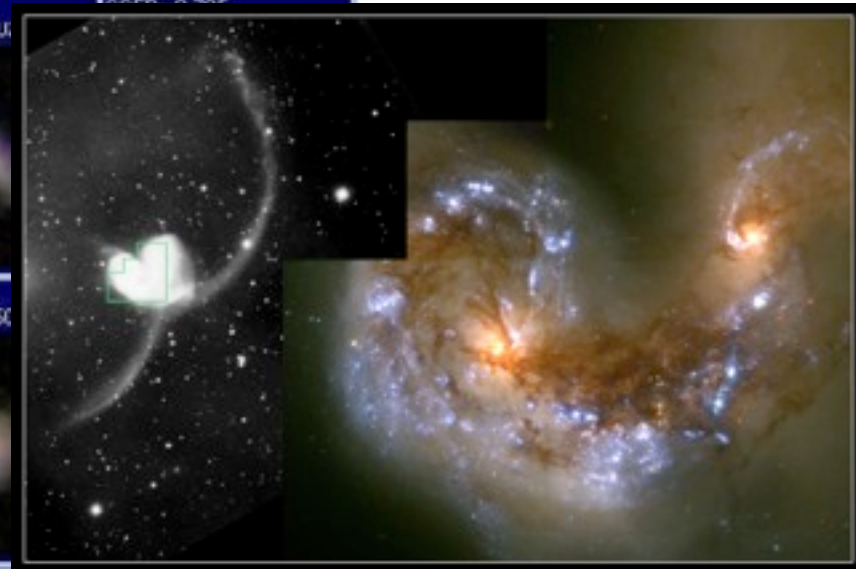
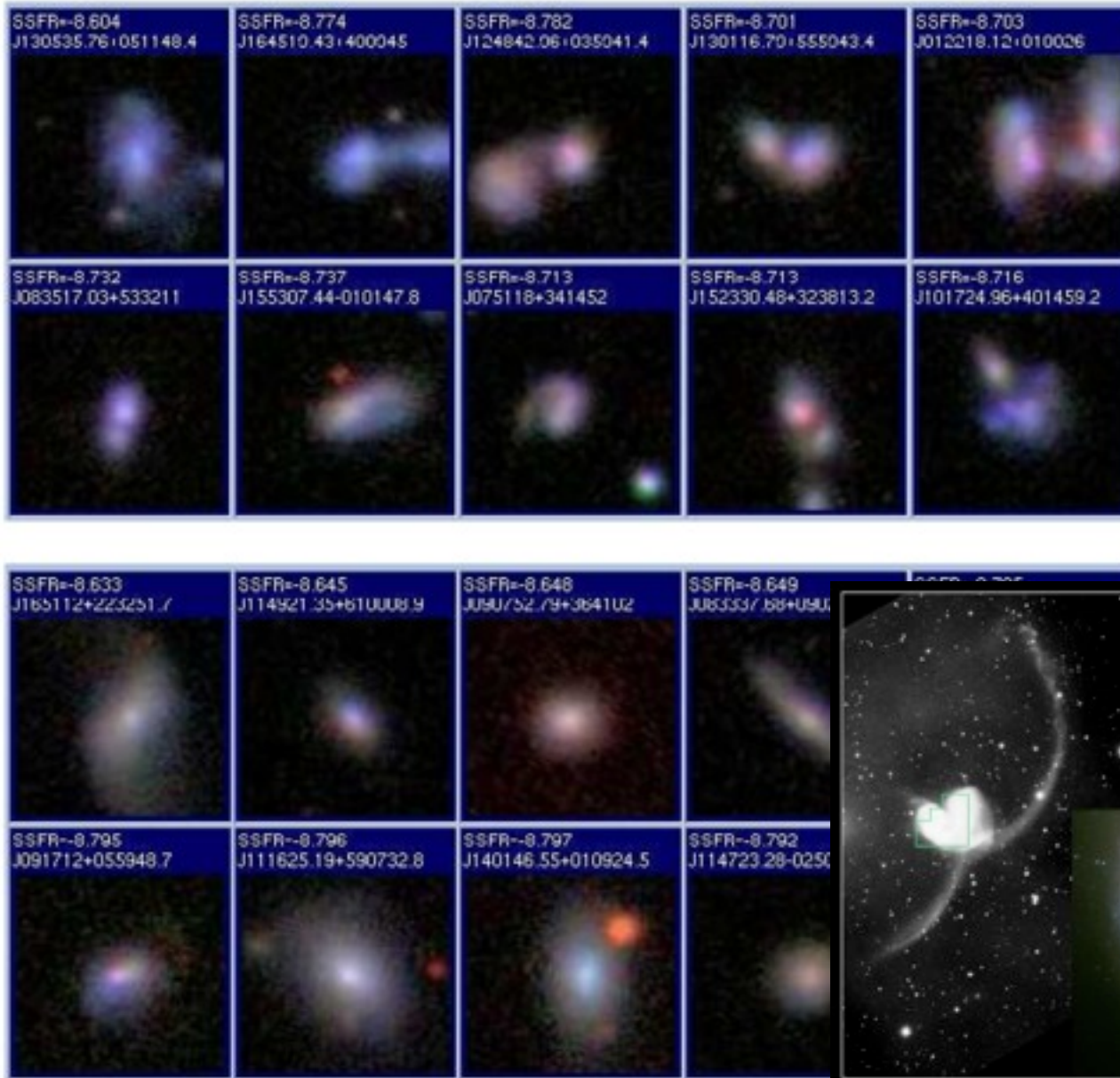


Merger rate from the close pair counts

Yipeng Jing
Shanghai JiaoTong University
Dept. of Phys. & Astron.

Collaborators:
Chunyan Jiang (SHAO), J. Han
(ICC, Durham)

Merging of galaxies



Merger of galaxies

- Observations: mergers are common and are important for understanding
 - Starbursts;
 - formation and evolution of elliptical galaxies, such as size evolution and mass growth;
 - maybe AGNs;
 - etc

How many mergers?

- Observations: galaxy image perturbed (substructures, irregular shape, etc) or the count of close pairs can be determined;
- Important: the timescale
 - of the substructure survival
 - of the merger of the satellite to the host

Theories

- Time scale can be only from theories. For the merger time:
 - Mechanism: the dynamical friction (Chandrasekhar);
 - Difficulties for analytical modeling
 - DM mass loss of the satellite;
 - dynamically evolution of the host halo;
 - motion of the central galaxy

- We employed a parallel version of the SPH code GADGET 2 (Springel 2005). The box is $100h^{-1}M_{pc}$ on a side, with 512^3 dark matter particles and 512^3 gas particles. Gravity is softened with a spline, roughly equivalent to a Plummer force softening of $4.9h^{-1}$ comoving kpc. There are totally 177 snapshots from $z=19$, among which 28 are before $z=3.5$, and 149 are at $z \leq 3.5$.

present time $z = 0$ with an equal logarithmic scale factor interval $\Delta \ln a = 0.01$ between two consecutive outputs. The large number of the outputs enables us to accurately sample orbits of satellites within massive haloes, with about 8 outputs for one dynamical crossing time. Both the good force resolution and the dense sampling of snapshots are crucial for the current study.

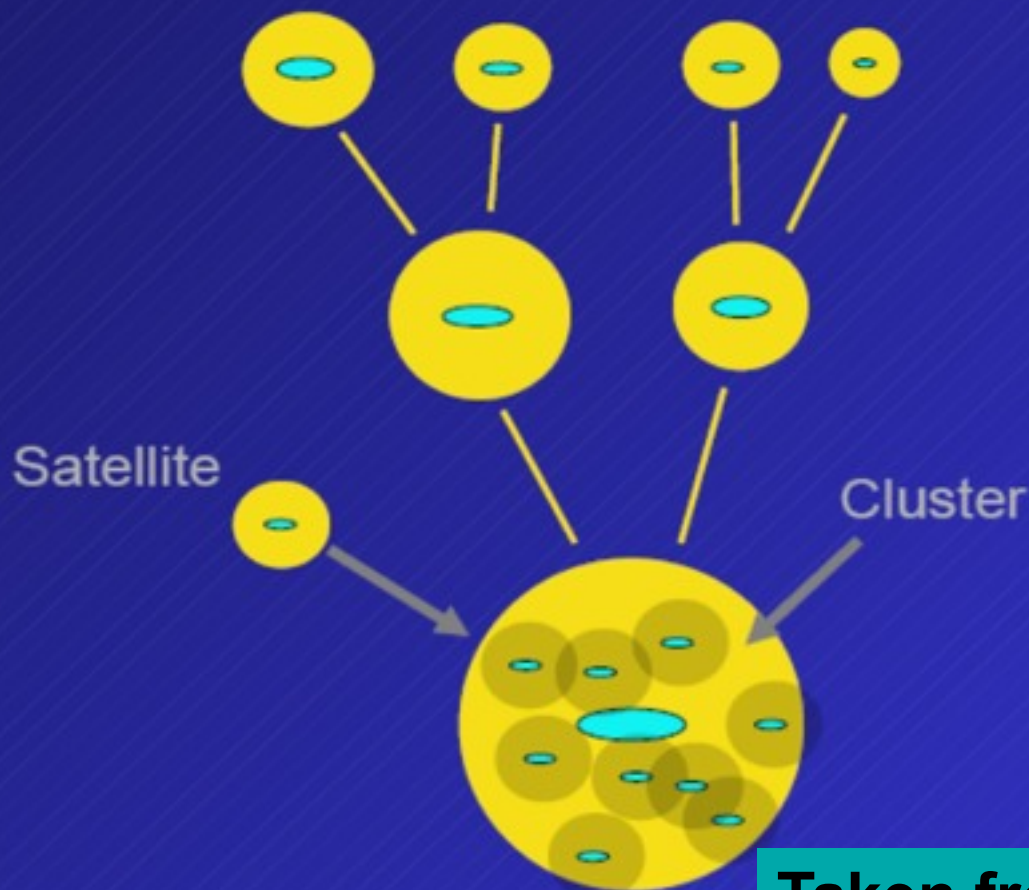
Two types of merger timescales in literature

- The time duration for a satellite falling into the central galaxy from the first crossing of the virial radius of host DM halo; important of theoretical modeling, such as in SAMs;
- The time duration for a close pair of galaxies at a fixed separation (small) to merge; important for observations

Theoretical framework for understanding evolution of galaxies and dark matter halos

A third key component: satellites vs. centrals

Smaller satellite galaxies can orbit for a time within larger halos without merging onto the central galaxies.



Taken from S. Faber

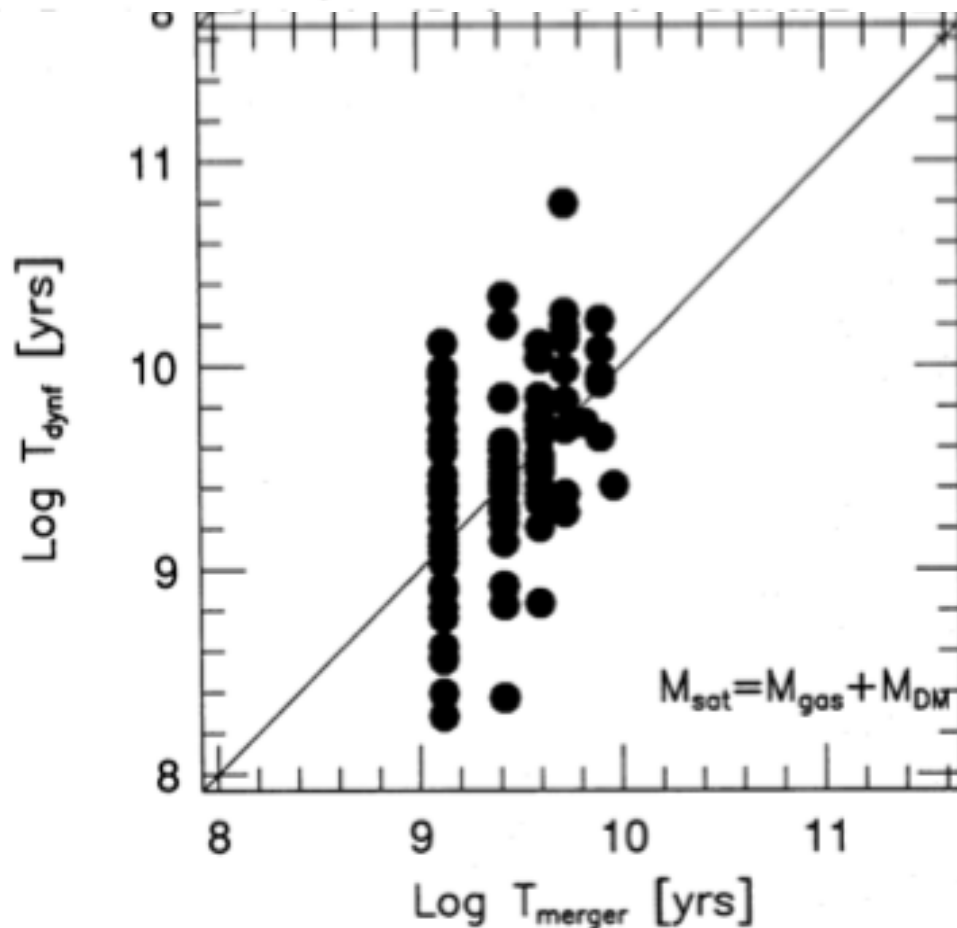
Taking into account the dependence on the orbital circularity, Lacey & Cole (1993) derived the following expression for the merger timescale of a satellite galaxy orbiting around a massive halo of circular velocity V_c ,

$$T_{\text{Chandra}} = 0.5 \frac{f(\epsilon) V_c r_c^2}{C G m_{\text{sat}} \ln \Lambda} \quad (1)$$

where ϵ is the circularity parameter of the satellite's orbit and r_c is the radius of a circular orbit with the same energy as the satellite's orbit. $f(\epsilon)$ describes the dependence of T_{Chandra} on the orbital circularity, and is approximated by $f(\epsilon) \sim \epsilon^{0.78}$ for $\epsilon > 0.02$ (Lacey & Cole 1993). C is a constant, approximately equal to 0.43, and m_{sat} is the satellite mass. $\ln \Lambda$ is the Coulomb logarithm, and is expected to be valid for the case

The assembly of galaxies in a hierarchically clustering universe

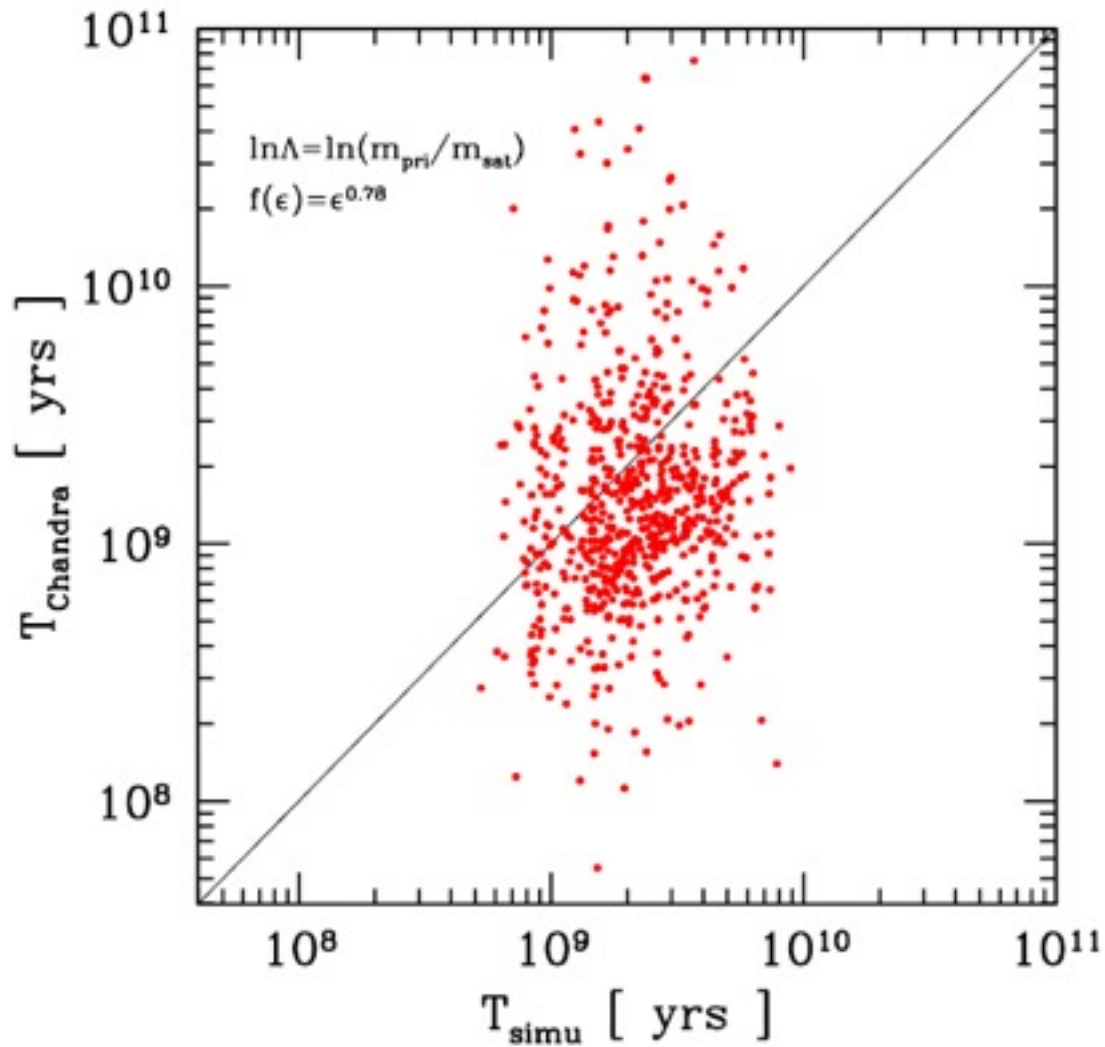
Julio F. Navarro,^{1*} Carlos S. Frenk¹ and Simon D. M. White^{2†}



$$\Lambda \equiv \frac{d_{\text{max}} V_{\text{typ}}^2}{G(m_{\text{sat}} + m_{\text{dm}})} = \frac{m_{\text{pri}}}{m_{\text{sat}}},$$

Using $\ln(m_h/m_{sub})$ for Coulomb logarithm (c.f. Navarro et al. (1995))

Predicted friction timescale

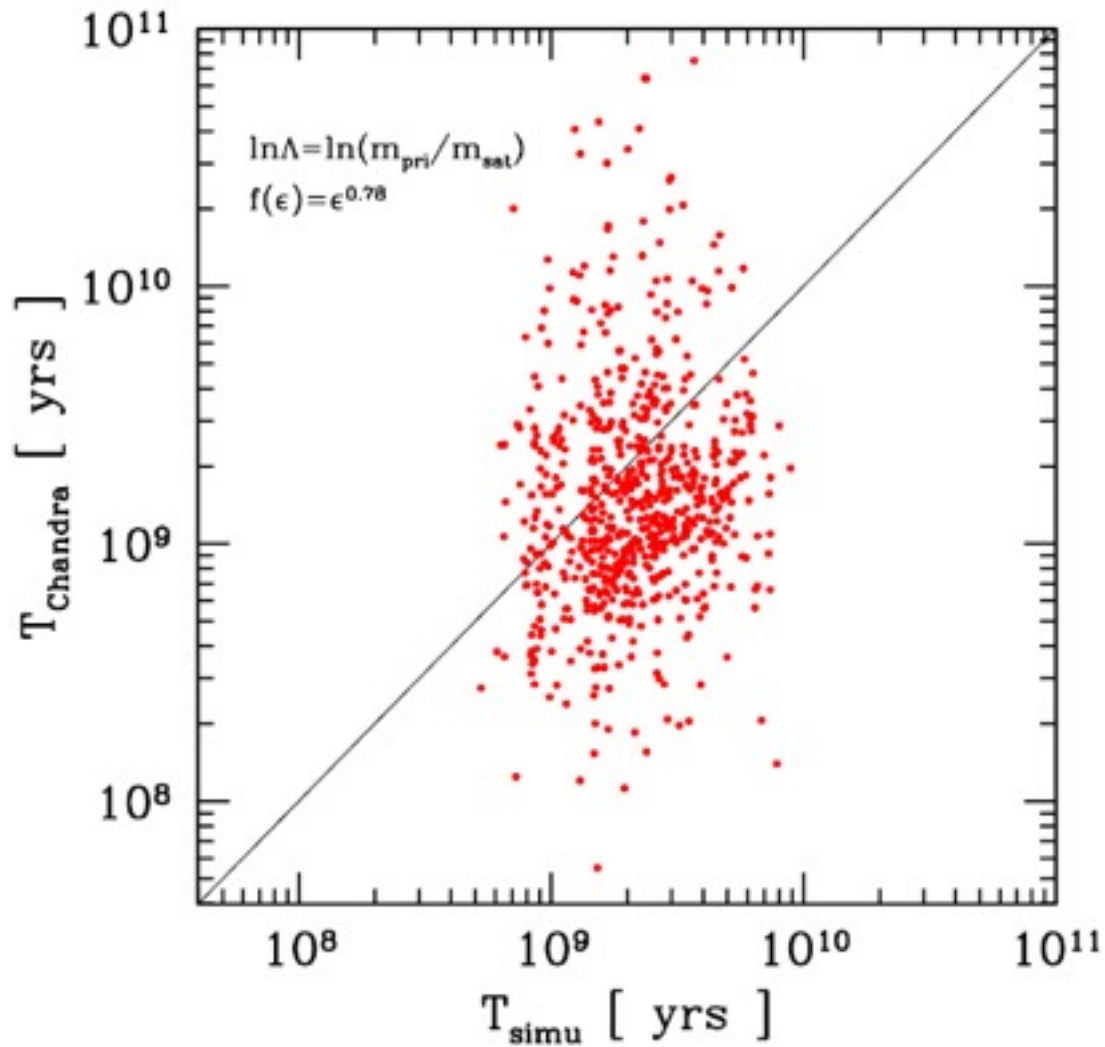


Measured merger timescale

Jiang et al. 2008

Using $\ln(m_h/m_{sub})$ for Coulomb logarithm (c.f. Navarro et al. (1995))

Predicted friction timescale



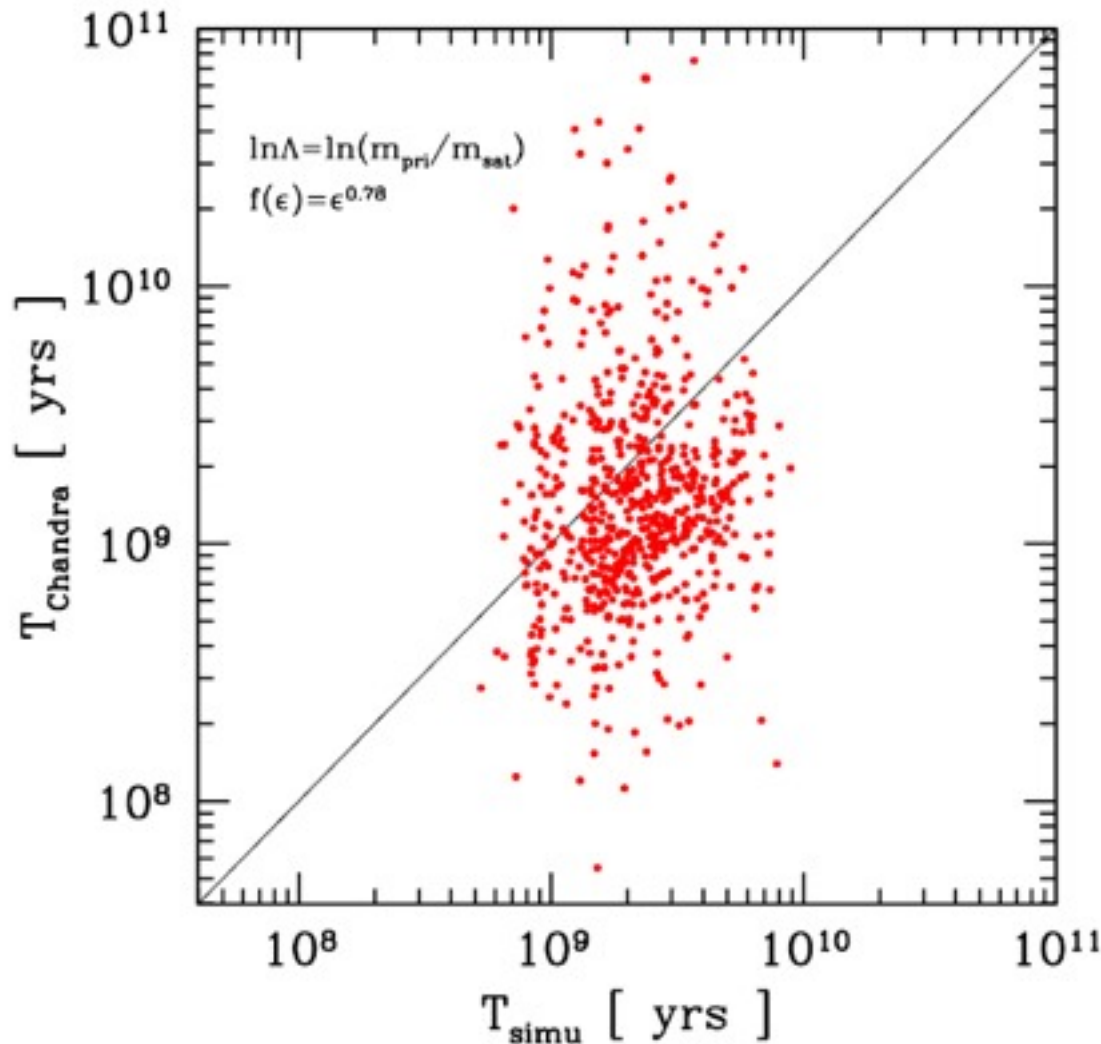
repeated
Navarro et
al. 1995

Measured merger timescale

Jiang et al. 2008

Using $\ln(m_h/m_{sub})$ for Coulomb logarithm (c.f. Navarro et al. (1995))

Predicted friction timescale



repeated
Navarro et
al. 1995

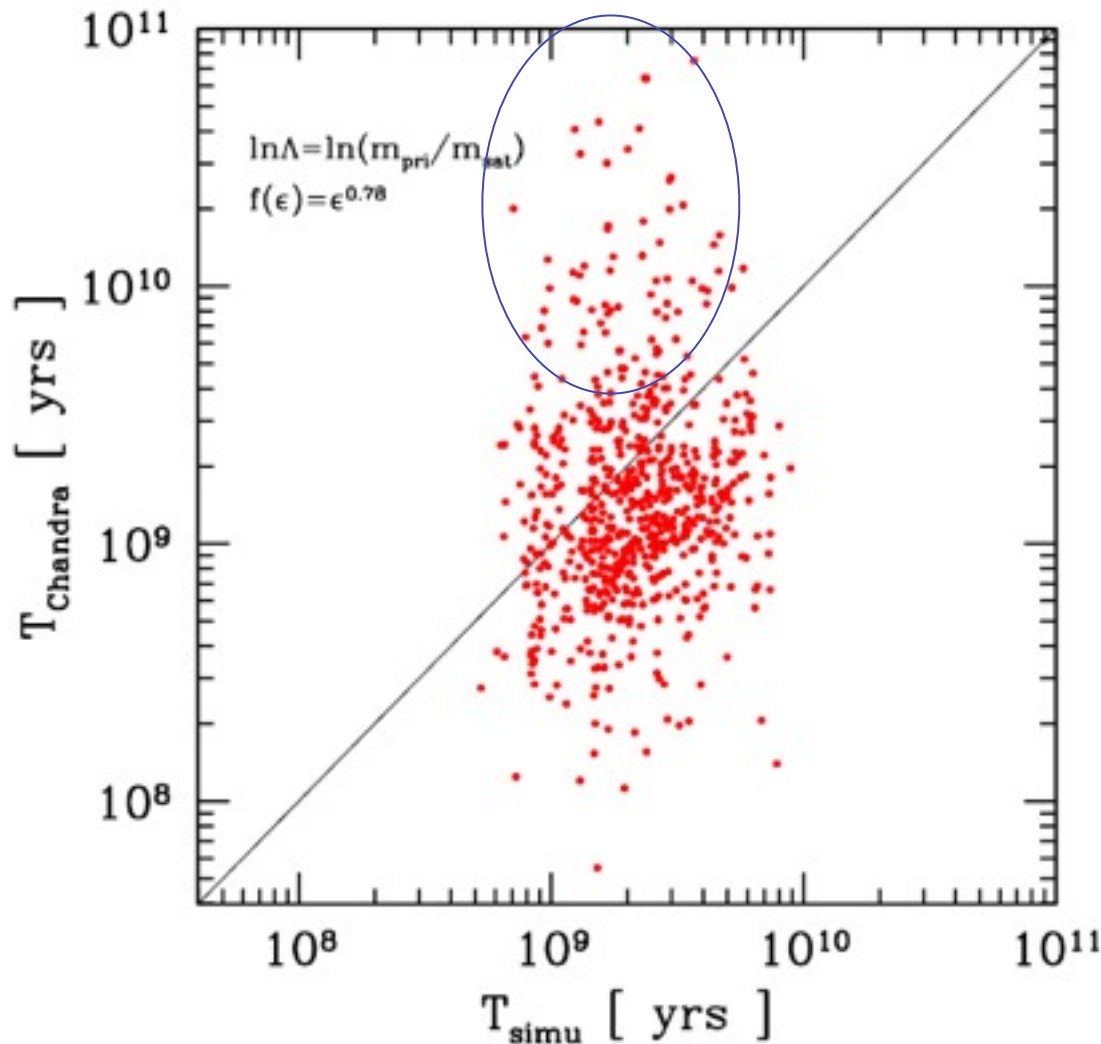
But our
conclusion
is different
from theirs

Measured merger timescale

Jiang et al. 2008

Using $\ln(m_h/m_{sub})$ for Coulomb logarithm (c.f. Navarro et al. (1995))

Predicted friction timescale



repeated
Navarro et
al. 1995

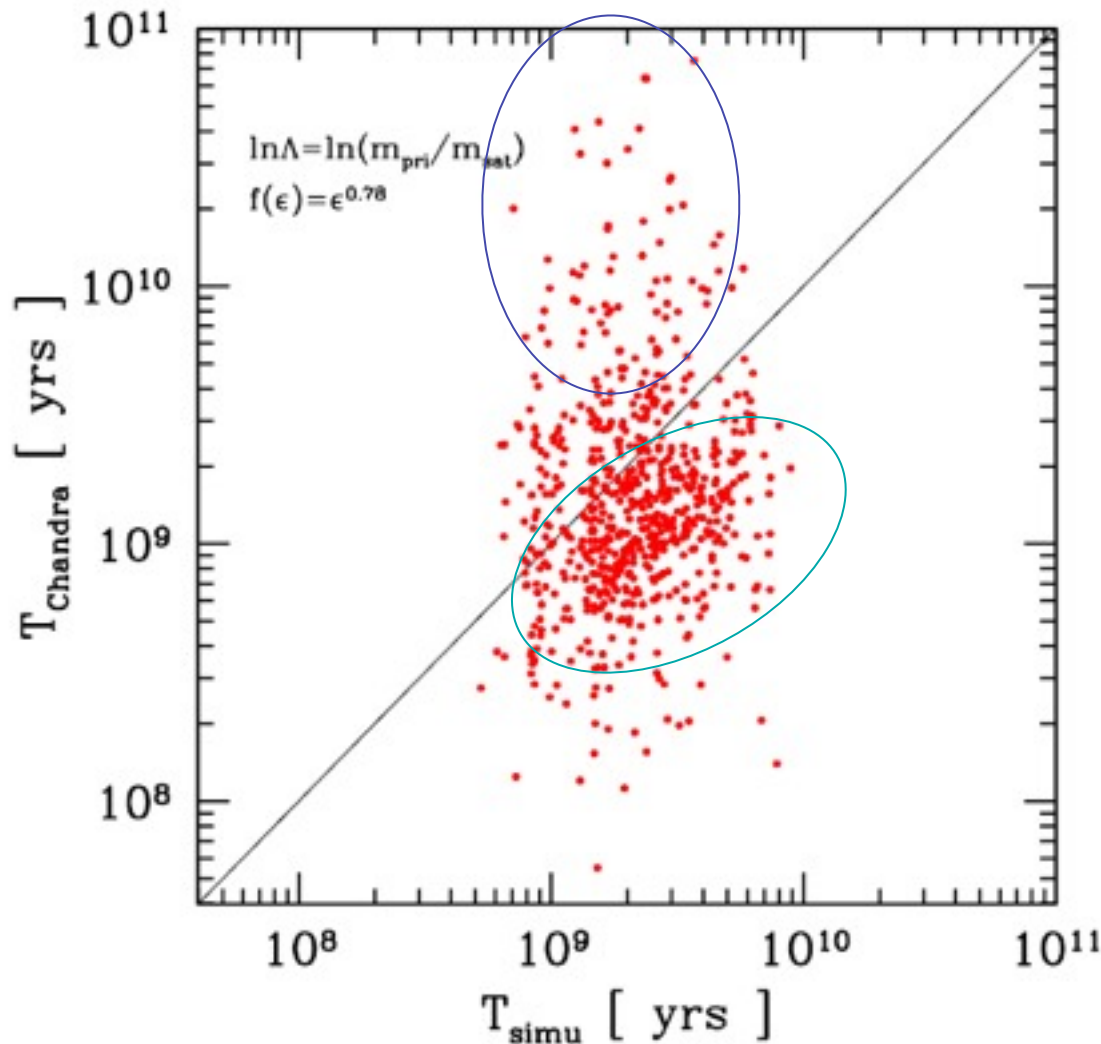
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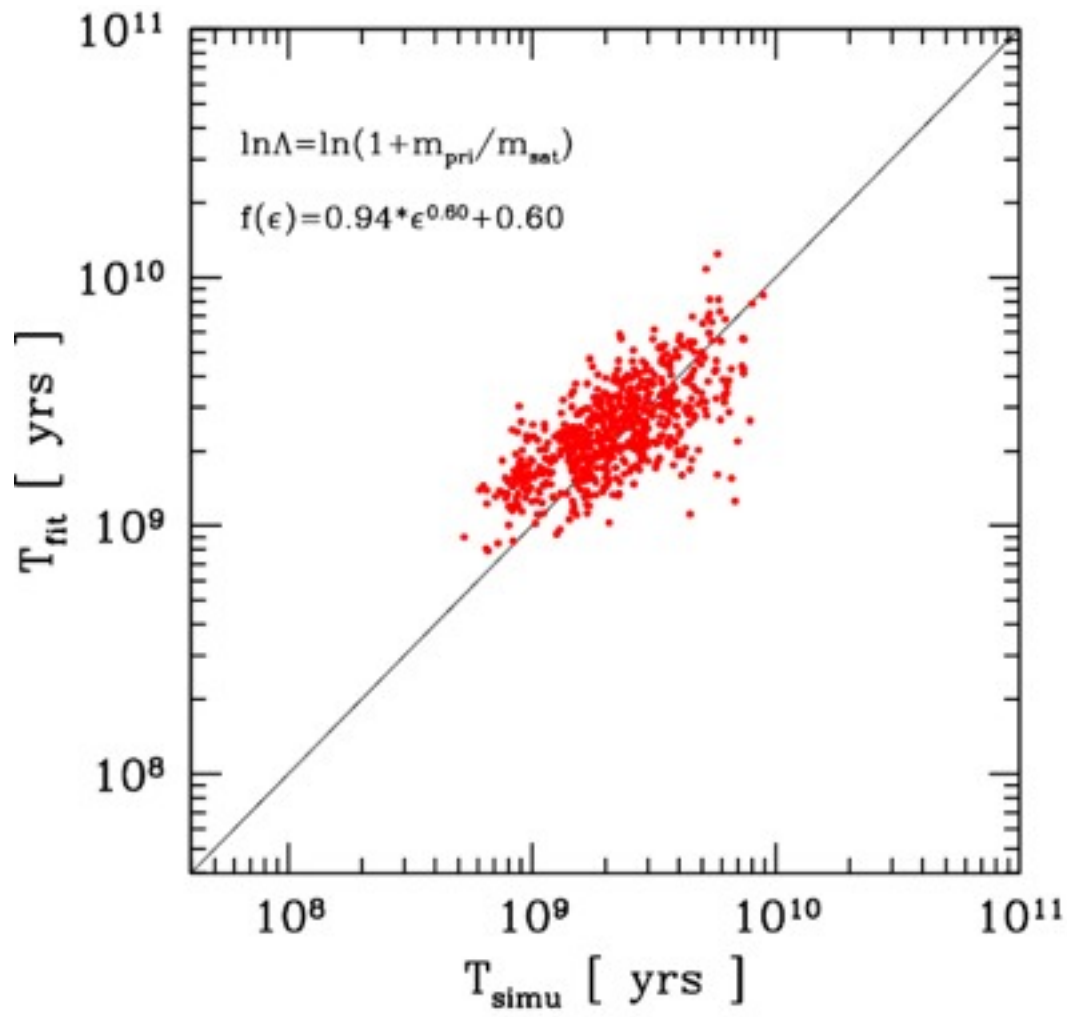
But our
conclusion
is different
from theirs

Measured merger timescale

Jiang et al. 2008

$$T_{\text{fit}} = \frac{0.90\epsilon^{0.47} + 0.60}{2C} \frac{m_{\text{pri}}}{m_{\text{sat}}} \frac{1}{\ln\left[1 + \left(\frac{m_{\text{pri}}}{m_{\text{sat}}}\right)\right]} \frac{\sqrt{r_{\text{vir}} r_c}}{V_c}, \quad (8)$$

Fitting formula



- Corrections:
- 1) **Mass loss**: a factor of 2 longer
 - 2) **Motion of the central galaxies and dynamical evolution**: weak dependence on ϵ (orbital circularity)
 - 3) **Dependence on the DM mass of the primary and satellite**: the Coulomb logarithm
 - 4) **Scatter**: 40% reflecting hierarchical formation and diversity of host halos

Merger timescale for all mergers

The Lives of High Redshift Mergers

Tom McCavana^{1*}, Miroslav Micic^{1,2}, Geraint F. Lewis¹, Manodeep Sinha³,
Sanjib Sharma¹, Kelly Holley-Bockelmann³ and Joss Bland-Hawthorn¹

¹*Sydney Institute for Astronomy, School of Physics, A28, The University of Sydney, NSW 2006, Australia*

²*Astronomical Observatory Belgrade, Volgina 7, P.O.Box 74 11060 Belgrade, Serbia*

³*Department of Physics & Astronomy, Vanderbilt University, Nashville, TN, USA*

2 May 2012

ABSTRACT

We present a comparative study of recent works on merger-timescales with dynamical friction and find a strong contrast between idealized/isolated mergers (Boylan-Kolchin et al. 2008) and mergers from a cosmological volume (Jiang et al. 2008). Our study measures the duration of mergers in a cosmological N-body simulation of dark matter, with emphasis on higher redshifts ($z \leq 10$) and a lower mass range. In our analysis we consider and compare two merger definitions; tidal disruption and coalescence. We find that the merger-time formula proposed by Jiang et al. (2008) describes our results well and conclude that cosmologically motivated merger-time formulae provide a more versatile and statistically robust approximation for practical applications such as semi-analytic/hybrid models.

arXiv:1204.6319

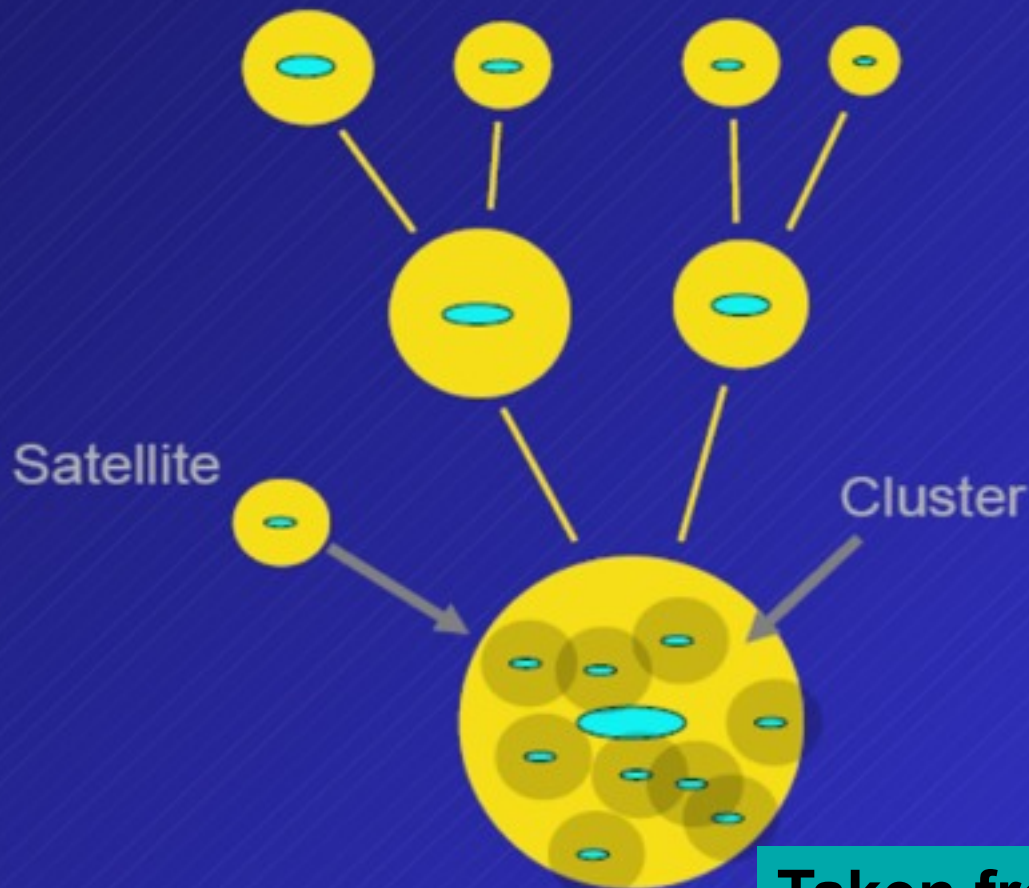
The second merger timescale

- A merger time for close pairs of certain mass (luminosity) and separation, related to measure the merger rate from the counts of close pairs in observations
- (Jiang, YPJ, Han, 2013, [astroph/1307.3322](#))

Theoretical framework for understanding evolution of galaxies and dark matter halos

A third key component: satellites vs. centrals

Smaller satellite galaxies can orbit for a time within larger halos without merging onto the central galaxies.



Taken from S. Faber

Some scaling considerations

$$T_{\text{fit}} = \frac{0.90\epsilon^{0.47} + 0.60}{2C} \frac{m_{\text{pri}}}{m_{\text{sat}}} \frac{1}{\ln\left[1 + \left(\frac{m_{\text{pri}}}{m_{\text{sat}}}\right)\right]} \frac{\sqrt{r_{\text{vir}} r_c}}{V_c}, \quad (8)$$

Considering $v_c \approx \sqrt{\frac{Gm_{1,v}}{r_{1,v}}}$ in the primary halo,

$$T_{\text{mg}} \propto \frac{m_{1,v}^{1/2} r_p^2}{G^{1/2} m_2 \ln \Lambda r_{1,v}^{1/2}}.$$

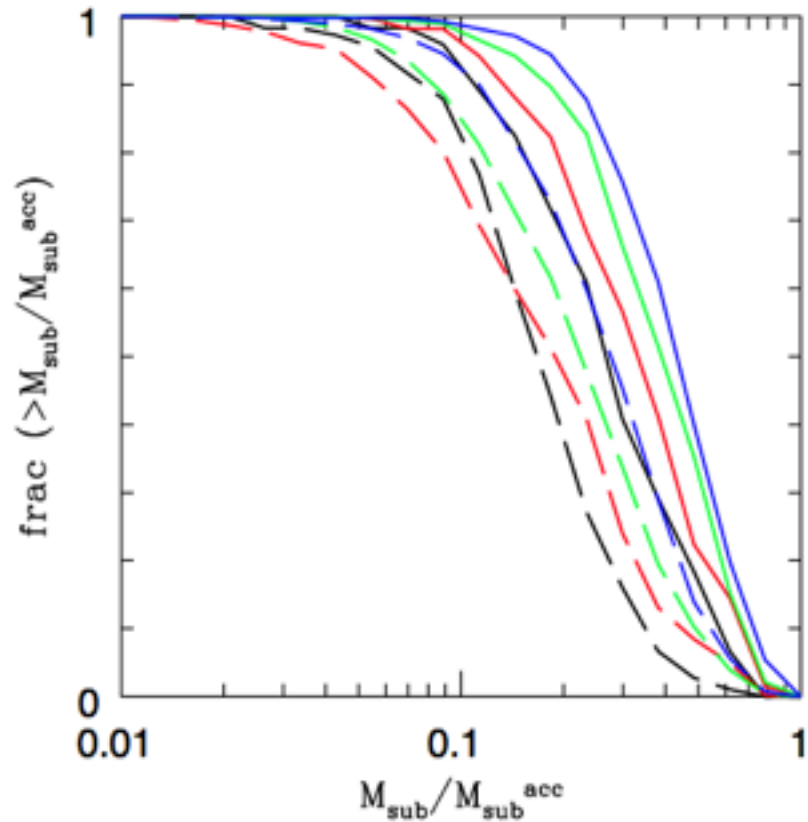
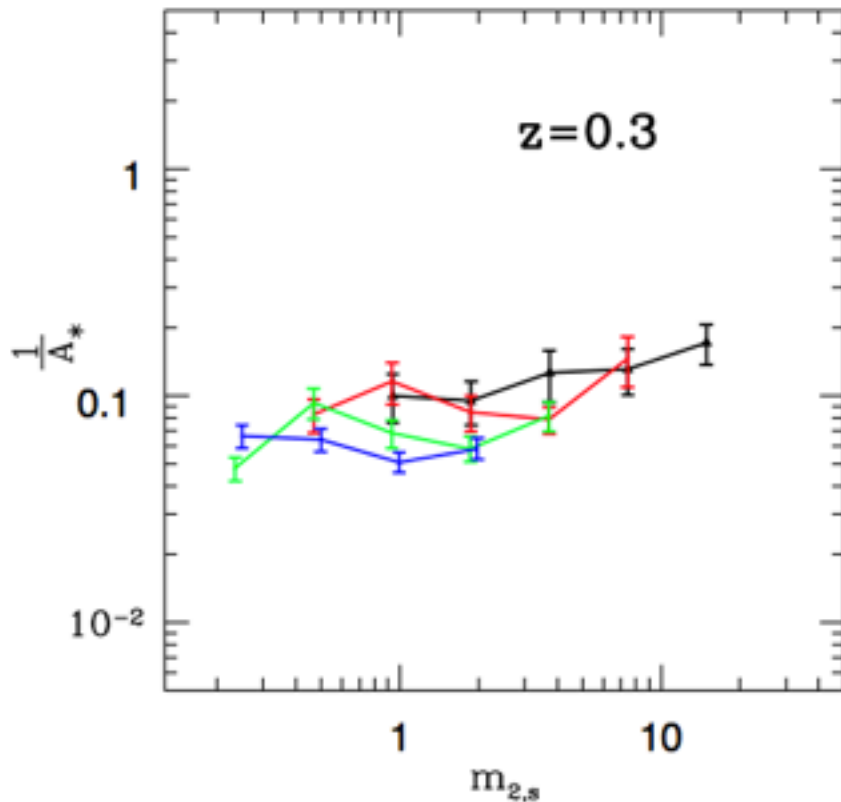
The volume merger rate can be written as

$$\Phi = C_{\text{mg}} n_1 n_p(< r_p) / T_{\text{mg}},$$

Replacing T_{mg} in equation (1) with equation (2), obtain

$$\Phi = A_* \frac{G^{1/2} m_2 r_{1,v}^{1/2} n_1 n_p(< r_p)}{m_{1,v}^{1/2} r_p^2}.$$

Stellar mass used for m_2



Distribution of the satellites

- 1) $r_p=77$ kpc/h (physical)
- 2) central stellar masses of 30.0 (black), 15.0 (red), 7.5 (green) and 4.0 (blue), respectively, in unit of $10^{10} h^{-1} M_{\odot}$
- 3) The DM retained by satellites is much larger than the stellar mass (at 77 kpc/h or 10-65 kpc/h dashed)

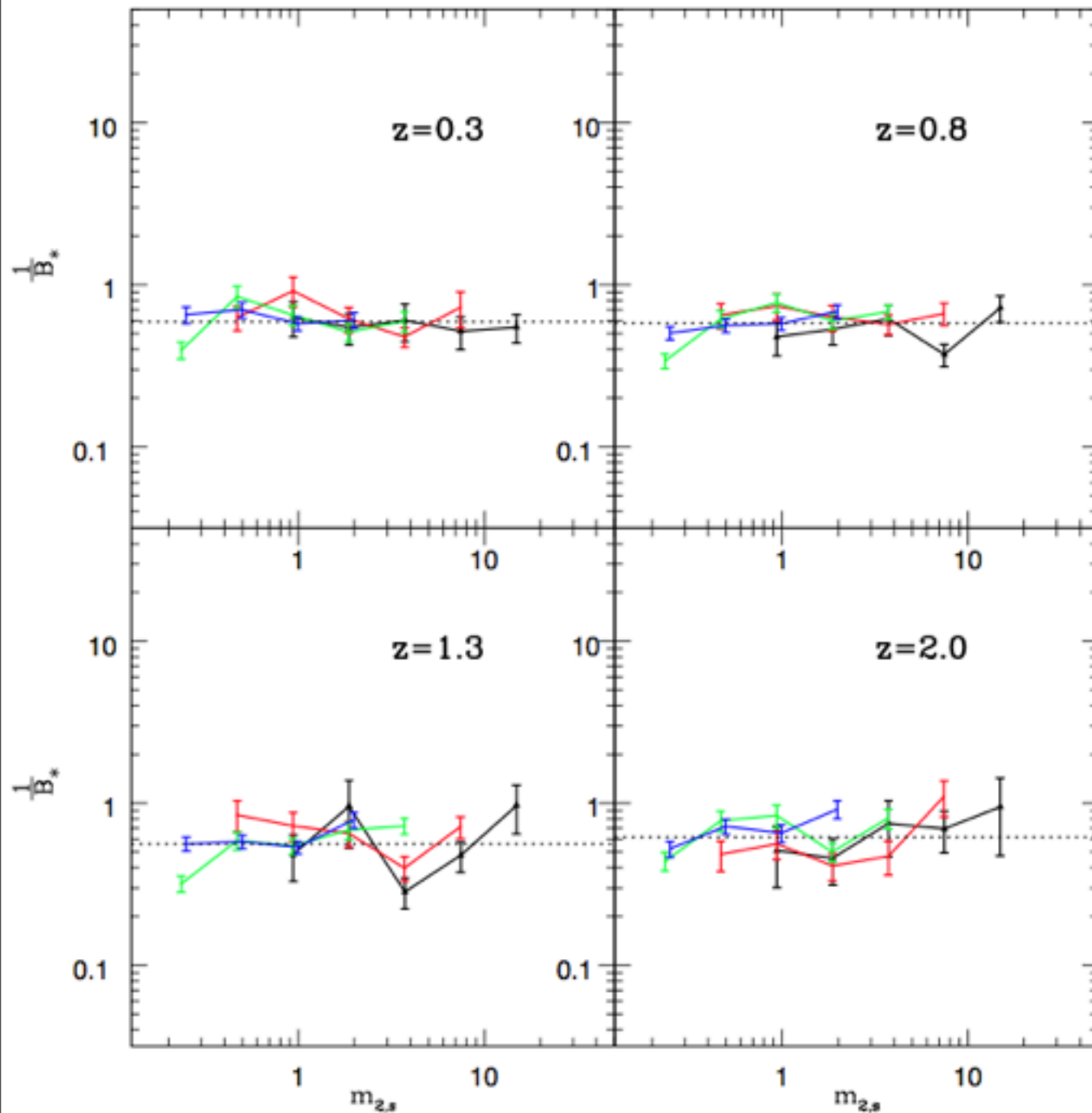
Retained mass ratio

More scaling considerations

- The retained mass of the satellite: $m_{2,v} \frac{r_p}{r_{1,v}}$
 - Correct when DM halo is an isothermal sphere of; but good for real DM halos
- With the definition of halos (200 critical density) and cosmological parameter relations, we have

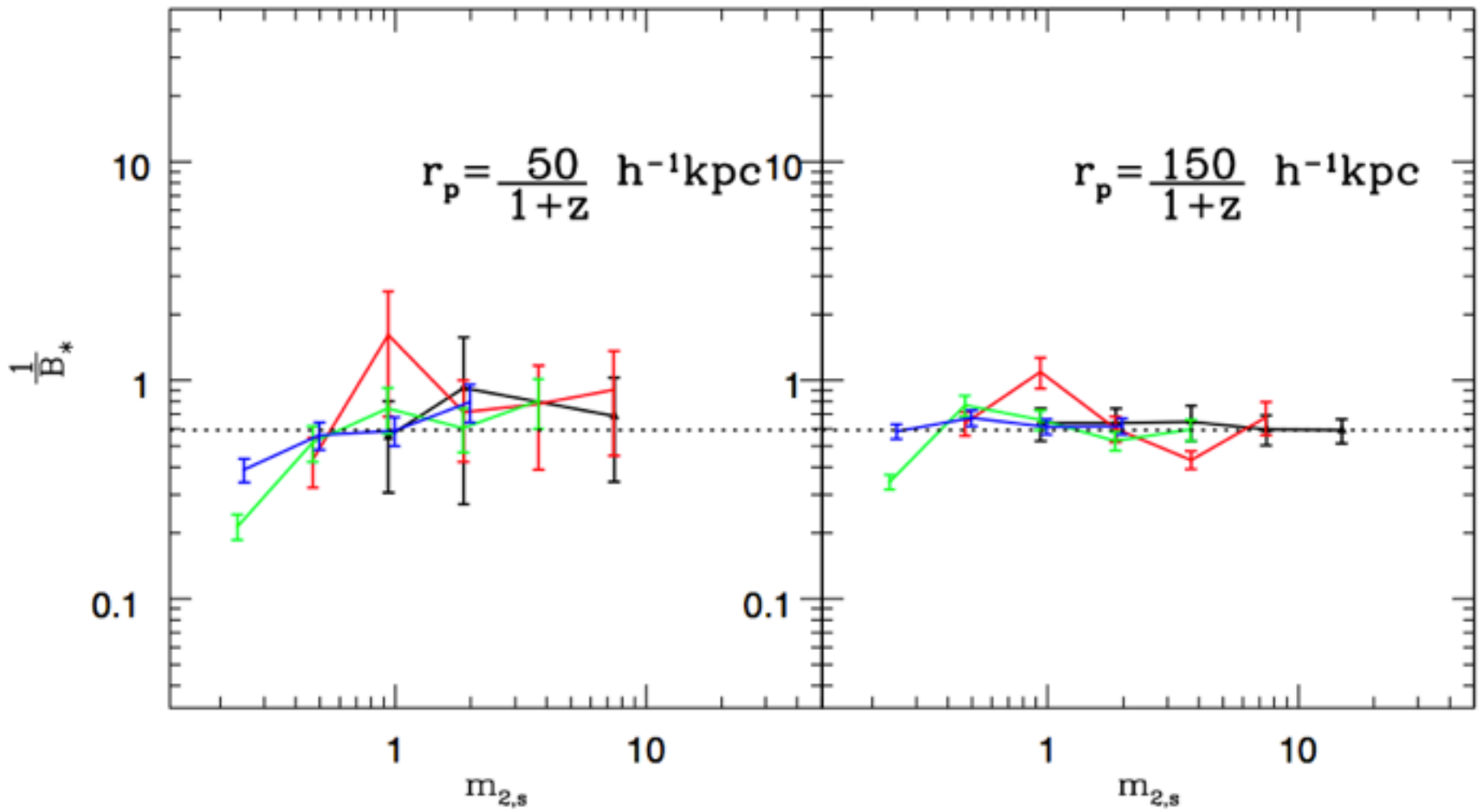
$$T_{\text{mg}} \propto \frac{m_{1,v}}{m_{2,v}} [m_{1,v} G H_0 E(z)]^{-1/3} r_p$$
$$\Phi = B_* \frac{m_{2,v} n_1 n_p (< r_p) [m_{1,v} G H_0 E(z)]^{1/3}}{m_{1,v} r_p}.$$

$E(z) = \Omega_\Lambda + \Omega_m(1+z)^3$: dimensionless Hubble parameter
(i.e. $H(z)$ in unit of H_0)

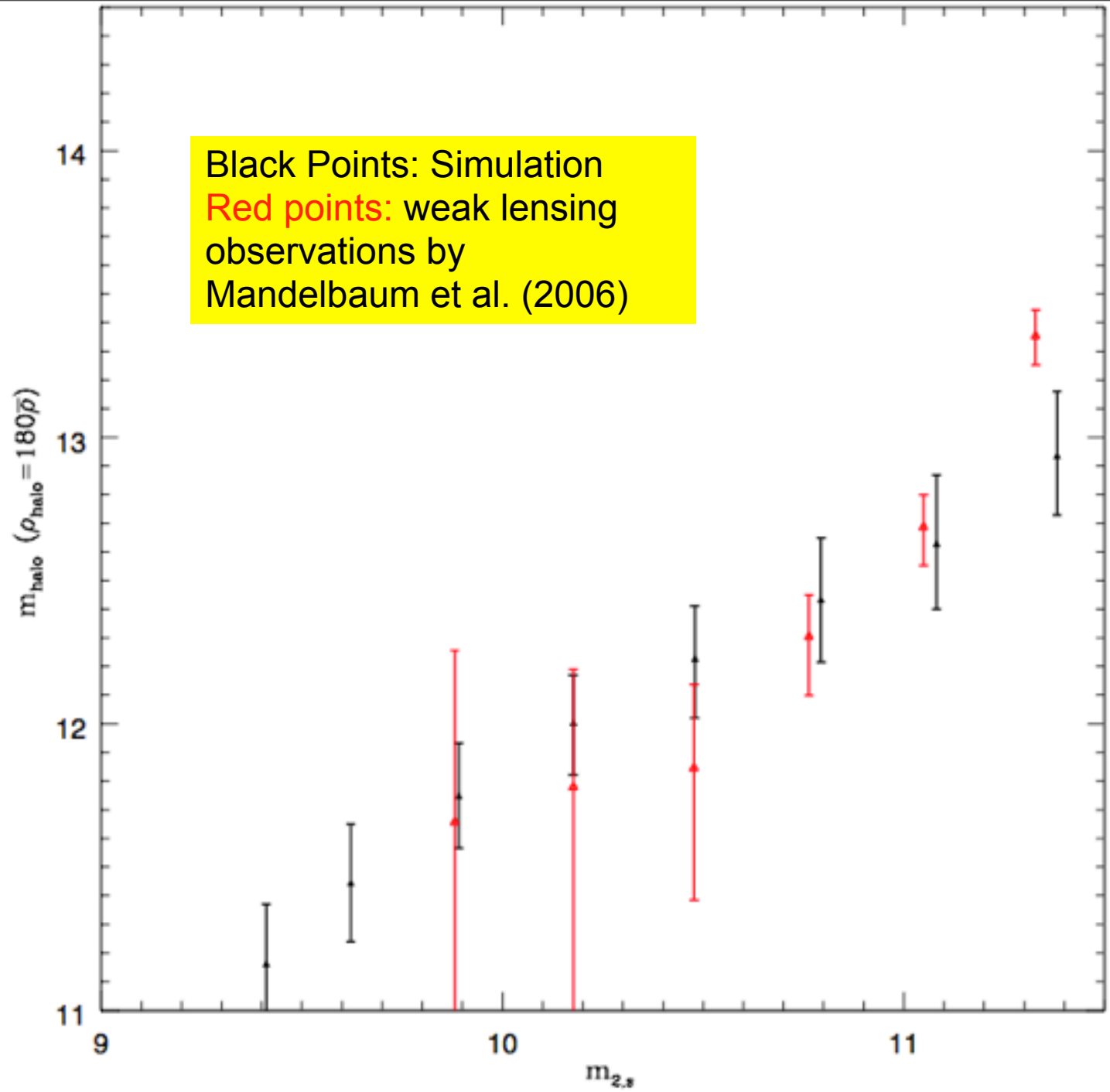


When the retained mass is considered for the satellites:

- 1) For different masses of central and satellites
- 2) For different redshifts



for the different separations



Applications to observations

- Measure the pair count per unit volume of stellar masses $m_{1,s}$ and $m_{2,s}$ (or luminosities) $N_p(< r_p) = n_1 n_p(< r_p)$

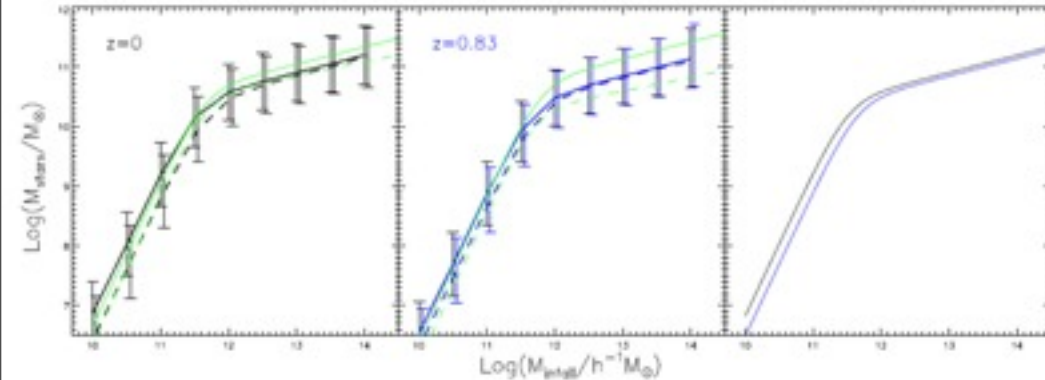
– n_1 is the density of galaxy 1 and n_p is the number count within projected r_p (corrected for the background) of galaxies 2 around galaxy 1

- Volume merger rate: $\Phi = N_p(< r_p) / T_{mg}$;

• Me

$$T_{mg}(< r_p^{\text{proj}}) = \frac{10^{-0.23}}{0.66} \frac{m_{1,v}}{m_{2,v}} [m_{1,v} G H_0 E(z)]^{-1/3} r_p^g$$

To be applied to observations: find out the stellar mass(luminosity)-halo mass relation, compute T_{mg} , and then the volume merger rate



HOD modeling from clustering and MF (Wang & YPJ 2010; cf. Yang et al 2003; Zheng et al 2005; Wang and Kauffmann 2006; Moster et al. 2011 etc)

- Many ways:
- 1)HOD, CLF etc
 - 2)Weak lensing
 - 3)Satellite dynamics
 - 4)Hydrodynamcs simulations etc
 - 5).....

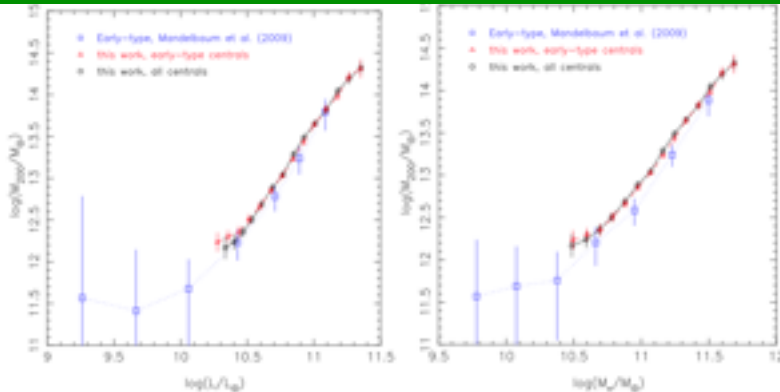
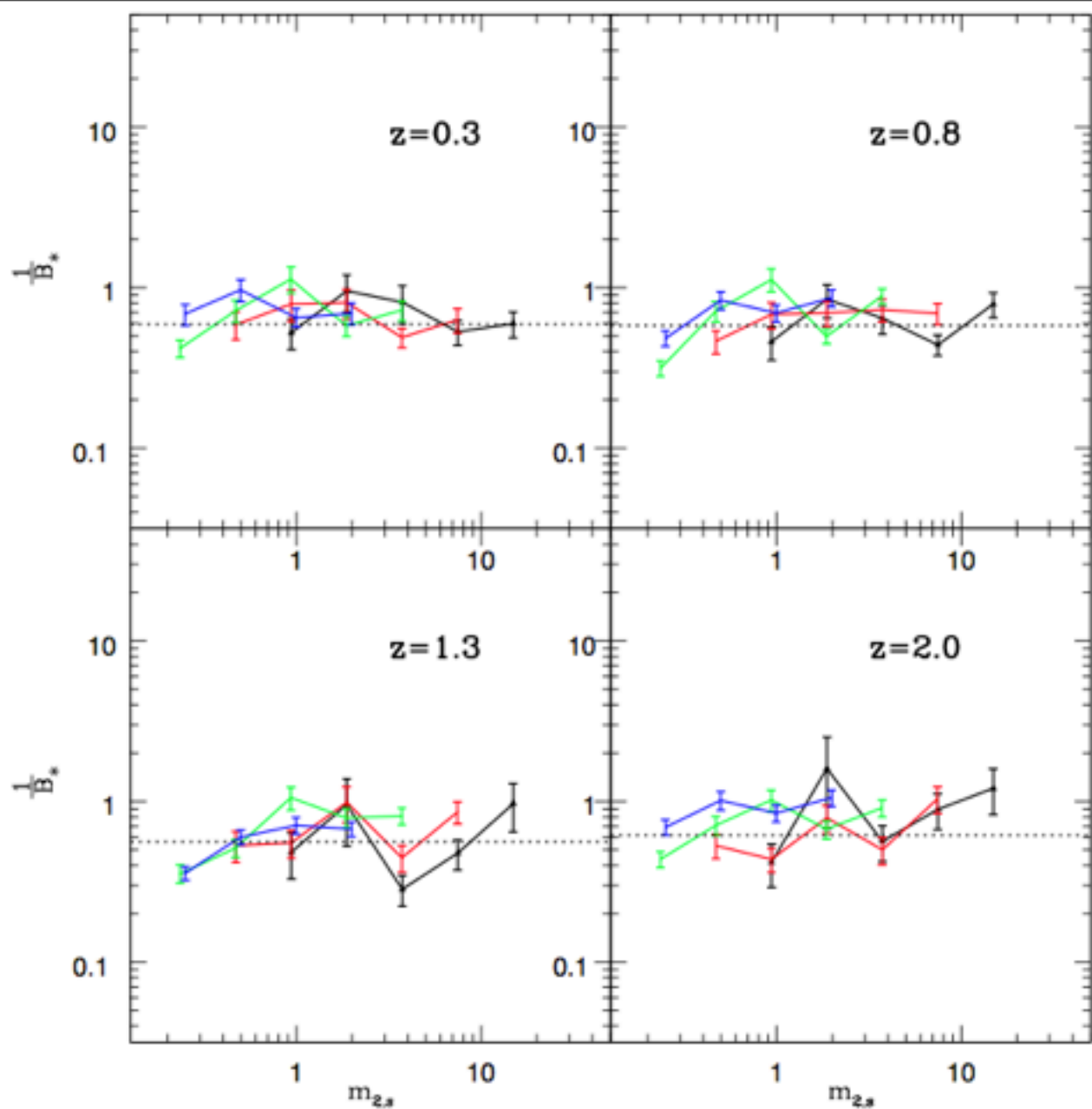


FIG. 18 — Dark matter halo mass as a function of galaxy luminosity. D&B-band (left) and stellar mass (right band (over)). Black circles

Direct measurement from weak lensing (Mandelbaum et al. 2006) and satellite dynamics (Li, et al 2012)

Conclusions and discussion

- We have considered the subhalo DM mass of satellites and the scaling relations on z and r_p ;
- We have, for the first time, **worked out a accurate merger timescale** that accounts for the dependences on the masses, redshift, and separation; **our results extends to minor mergers (1:30 in M_s)**
- The mass dependence is important since M_h - M_s could be steep [e.g Kitzbichler & White (2008) for major mergers (1:1-1:4) but valid for 1:4 mergers only];



When a different method to define a “galaxy” ($b=0.0125$), with the stellar mass much smaller than the fiducial case ($b=0.025$)